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The power delivered to any device as a function of time is given by the product of the instantaneous voltage across the device and the instantaneous current through it.

For an inductor the following is true:

$$p = v_L i_L = L i_L \frac{di}{dt} = \frac{1}{L} v_L \int_{-\infty}^t v(t) dt$$

For a single series R-L circuit fed by a voltage unit step  $V_o u(t)$  the following equations can be written:

$$-V_o u(t) + L \frac{di}{dt} + i R = 0 \quad \text{Using partial fraction decomposition, this expression for } I(s) \text{ can be written to facilitate determining the time expression.}$$

$$V_o u(t) = L \frac{di}{dt} + i R$$

$$\frac{V_o}{s} = I(s) [sL + R]$$

$$I(s) = \frac{V_o}{s[sL + R]} = \frac{A}{s} + \frac{B}{sL + R}$$

$$I(s) = \frac{V_o}{s[sL + R]} \quad A = \frac{V_o}{R} \quad B = -AL = -V_o \frac{L}{R}$$

Putting all these expressions together yields the final expression for the current

$$I(s) = \frac{V_o/R}{s} + \frac{-A/R}{s + R/L} = \frac{V_o}{R} \left[ 1 - e^{-\frac{Rt}{L}} \right] u(t)$$

The power delivered by the source or absorbed by the passive network is then:

$$p = vi = \frac{V_o^2}{R} \left[ 1 - e^{-\frac{Rt}{L}} \right]^2 u(t) \quad [1]$$

Power delivered to the resistor is

$$p_R = vi = i^2 R = \frac{V_o^2}{R} \left[ 1 - e^{-\frac{Rt}{L}} \right]^2 u(t) \quad [2]$$

The power delivered to the inductor is slightly more complex to determine.

$$v_L = L \frac{di}{dt} \quad \text{Taking the derivative of the current w.r.t. time gives:}$$

$$\frac{di(t)}{dt} = \frac{d}{dt} \left[ \left( \frac{V_o}{R} - \frac{V_o}{R} e^{-\frac{Rt}{L}} \right) u(t) \right] = \frac{V_o}{R} \delta(t) - \frac{V_o}{R} \left[ e^{-\frac{Rt}{L}} \delta(t) - \frac{R}{L} e^{-\frac{Rt}{L}} u(t) \right]$$

Combining things and noticing the cancellation of the Dirac delta terms at  $t = 0$  gives the following for the voltage across the inductor:

$$v_L = \frac{V_o}{L} e^{-\frac{Rt}{L}} u(t) \quad [3]$$

Now, for the power developed in the inductor, we multiply the time-domain voltage and current together to obtain

$$p_L = v_L i = L \frac{di}{dt} i = L \left[ \frac{V_o}{L} e^{-\frac{Rt}{L}} u(t) \right] \frac{V_o}{R} \left( 1 - e^{-\frac{Rt}{L}} \right) \quad [4, 5]$$

$$p_L = \frac{V_o^2}{R} e^{-\frac{Rt}{L}} \left( 1 - e^{-\frac{Rt}{L}} \right) u(t)$$

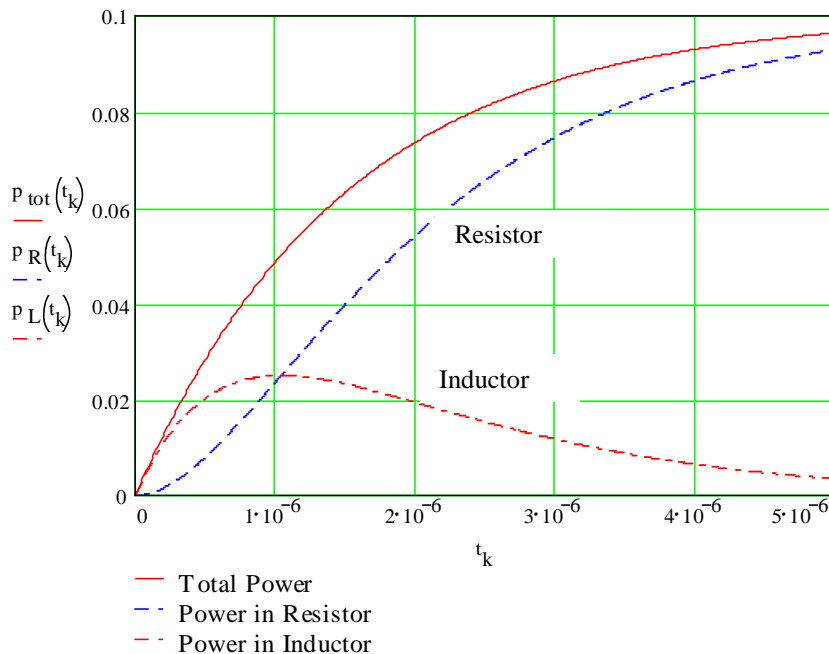


Figure 1

As time progresses the total power in the circuit increases up to the maximum possible of 0.1 Watt. This is set by the value of the resistor: 10 ohms. The L and C values used were 4.7  $\mu$ H and 15  $\mu$ F. One time constant for the values used is equal to 1.5  $\mu$ sec.

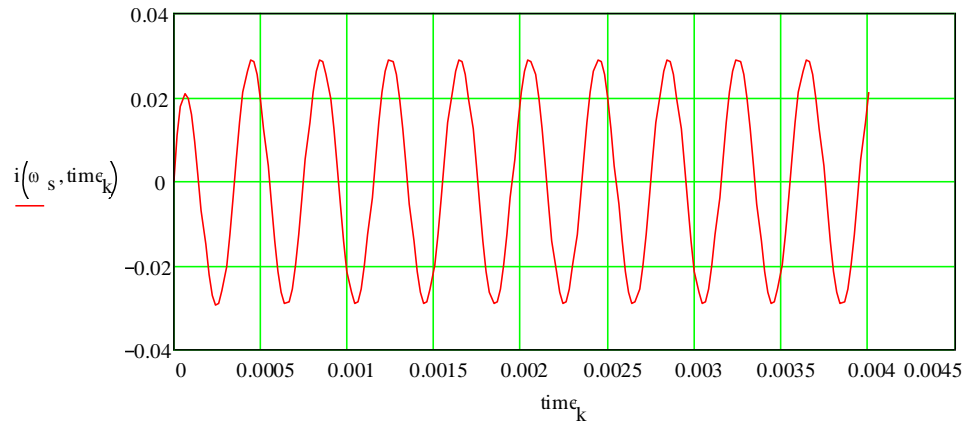
The input stimulus for this problem was a unit step function of magnitude 1.0 volts.

Let us now replace the input unit step function with a sinusoidal source of same magnitude  $V_o$ . Applying Laplace-type analysis gives us a similar expression to consider, however with increased complexity.

$$-V_o \cos(\omega_s t) + L \frac{di}{dt} + i(t) = 0 \quad \Rightarrow \quad \frac{V_o s}{s^2 + \omega_s^2} = I(s)(sL + R)$$

$$\frac{V_o s}{(s^2 + \omega_s^2)(sL + R)} = \frac{As + B}{s^2 + \omega_s^2} + \frac{C}{sL + R} \Rightarrow A = \frac{V_o R}{(\omega_s L)^2 + R^2} \quad B = \frac{V_o L \omega_s^2}{(\omega_s L)^2 + R^2} \quad C = \frac{L(L - RV_o)}{(\omega_s L)^2 + R^2}$$

$$I(s) = \frac{V_o s}{(s^2 + \omega_s^2)(sL + R)} \quad [6]$$



Taking the inverse Laplace transform of this last expression gives the current  $I(t)$  which is graphed as a function of time in Figure 2. As is seen in Figure 2, the steady-state current takes a small period of time to build up after the transient response.

$$i(t) = A \cos(\omega_s t) + \frac{B}{\omega_s} \sin(\omega_s t) + \frac{C}{L} e^{-\frac{Rt}{L}} \quad [7]$$

The power delivered to the resistor is straight forward with the expression for current in [7]:  $I(t)^2 R$ . Using the following expression for  $di/dt$ , the voltage across the inductor  $L$  can be determined, then multiplied by  $i(t)$  to determine the instantaneous power delivered to the inductor as well.

$$v_L = L \frac{di}{dt} = L \left[ -A\omega_s \sin(\omega_s t) + \frac{B\omega_s}{\omega_s} \cos(\omega_s t) + \frac{RC}{L^2} e^{-\frac{Rt}{L}} \right] \quad [8]$$

The power delivered to the inductor is then  $v_L(t) i(t)$ .

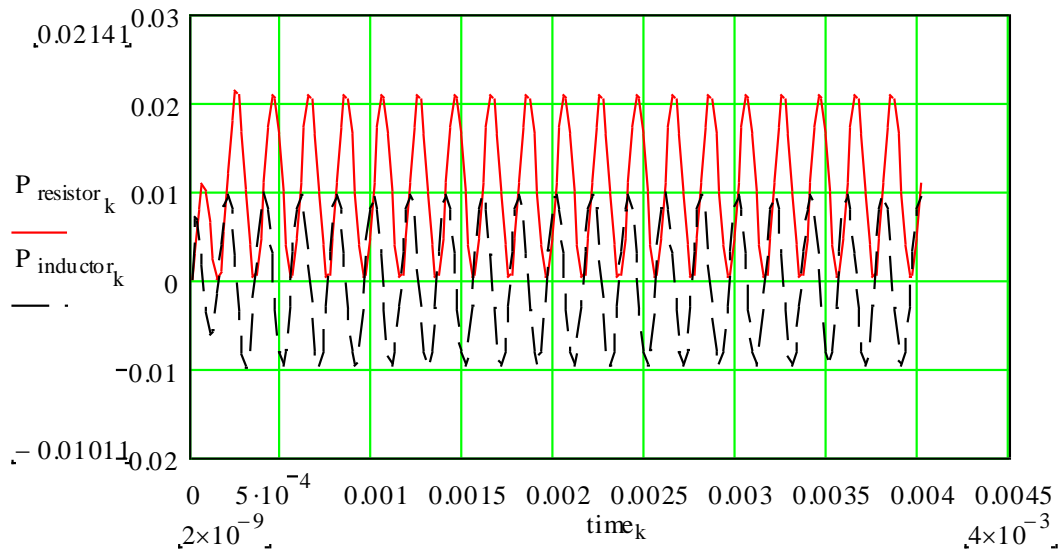


Figure 3

Adding the individual powers developed in the resistor and inductor gives the same total power delivered to the circuit, found by multiplying  $v(t)$  by  $i(t)$ .

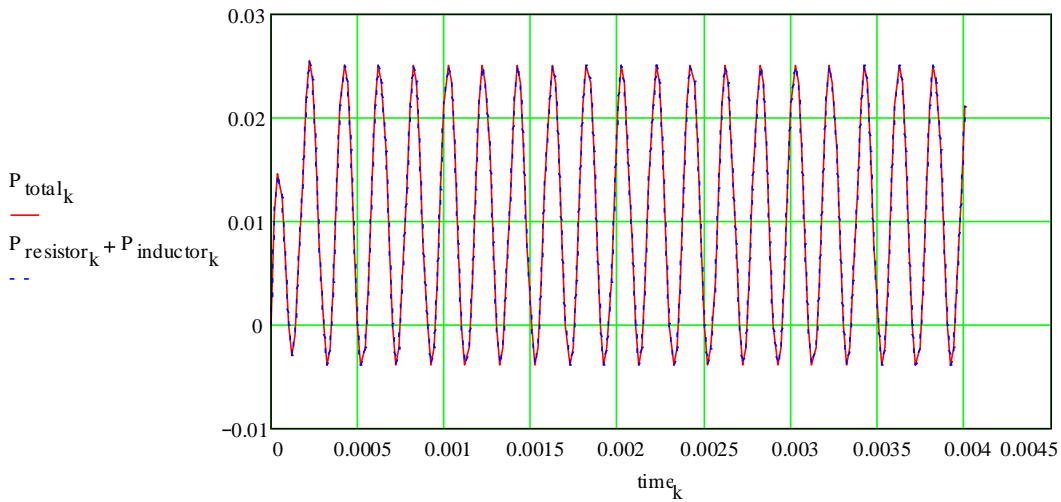
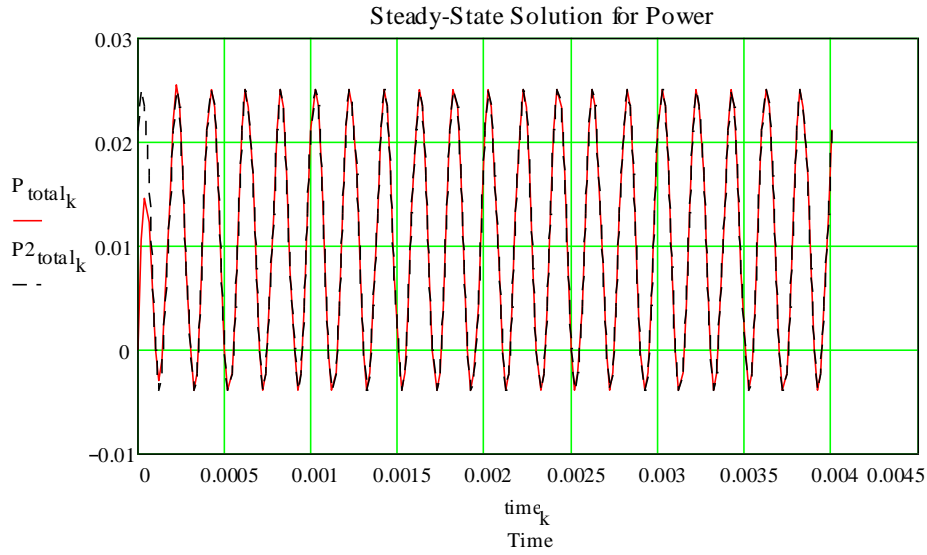


Figure 4

Total Power Equals the Summation of the Individual Powers



*Figure 5*  
Steady-State Power Lacks Initial Transient Response Found in the Complete Solution

